# Dynamics of Pulsar Wind Nebula with Shuta Tanka Magnetic Dissipation and Turbulence Aoyama Gakuin Univ. ABSTRACT

Pulsar wind nebulae (PWNe) are composed of relativistic magnetized plasma wind supplied from their central pulsars, called pulsar wind. Based on standard ideal magnetohydrodynamics model of PWN, magnetization of the pulsar wind is much less than unity in order to explain the observed slow expansion of the PWN. On the other hand, the pulsar wind is Poynting dominated at its base. This is called  $\sigma$ -problem. Here, we introduce our recent study extending the standard ideal MHD model by Kennel & Coroniti. Including the effects of magnetic dissipation and turbulence, requirement of low-magnetization from the dynamics of PWNe is relaxed.

 $\sigma \equiv \frac{\text{Poynting flux}}{\text{Plasma energy flux}}$ We extend KC model including the I.  $\sigma$ -problem in PWNe III. Model three non-ideal MHD effects. Tanaka, Toma&Tominaga(2018) Pulsar winds are a high- $\sigma$  flow at their origin, i.e., pulsar magnetosphere, while they seem to be a low- $\sigma$ **Dissipation of turbulent B-field** flow at their end, i.e., termination shock or PWN. ③ Radiative (synchrotron) loss of plasma heat

#### **σ-problem from dynamical properties**



Observed expansion velocity of the Crab Nebula  $v_{\rm Crab}$  of ~ 1,500 km/s is much slower than speed of light *c*.  $v_{\text{Crab}}$  /  $c \sim \sigma$  from the standard model by Kennel & Coroniti, i.e., we require  $\sigma \ll 1$ .

### **σ-problem from spectral properties**



**Observed broadband** spectrum of the Crab Nebula is explained by a emission from low-σ plasma (synchrotron + inverse Compton scattering model).

 $\langle \nabla_{\mu}(nu^{\mu}) \rangle = 0,$  $\langle \nabla_{\mu} T^{\mu t} \rangle = -\gamma \frac{\Lambda_{\text{rad}}}{M}, \quad \text{Conservation of total energy}$  $-\langle u_{\nu}\partial_{\mu}T_{\rm FL}^{\mu\nu}\rangle = \frac{\delta b^2/2}{\tau_{\rm diss}} - \frac{\Lambda_{\rm rad}}{c}, \quad \begin{array}{c} \text{Conservation of} \\ \text{fluid internal energy} \end{array}$  $\frac{1}{2} \langle \bar{b}_{\mu} e^{\mu\nu\alpha\beta} \nabla_{\nu} F_{\alpha\beta} \rangle = -\frac{\bar{b}^2/2}{\tau_{\text{conv}}}, \qquad \text{Induction equation for}$  $\frac{1}{2} \langle \delta b_{\mu} e^{\mu\nu\alpha\beta} \nabla_{\nu} F_{\alpha\beta} \rangle = \frac{\overline{b}^2/2}{\tau_{\text{conv}}} - \frac{\delta b^2/2}{\tau_{\text{diss}}}. \quad \text{(mean) toroidal field + turbulent field}$ 

Eq. of continuity

- Steady & spherical symmetry
- Relativistically hot plasma (Eq. of state)
- Five variables  $(n, h, p, u, b^2, \delta b^2)$
- Omitting all the RHS terms, KC model is recovered.
- $\Lambda_{\rm rad}$ ,  $\tau_{\rm diss}$ ,  $\tau_{\rm conv}$  are the parameters of the system.

#### We focus on dynamical $\sigma$ -problem here.

## II. Model of Kennel & Coronti(1984)

SNR G21.5-0.9 PWN G21.5-0.9 PSR J1833-1034 G21.5-0.9(Chandra)

- 1) Cold e<sup>±</sup> plasma from pulsar
- ② Standing termination shock 3 Steady and spherical system
- **4** Pure radial outflow
- **5** Pure toroidal B-field

**6** Ideal MHD approximation



We discard the above approx. (5) & (6).

V. Further Studies

Broadband brightness & polarization profiles of PWNe, (1) i.e., solving spectral  $\sigma$ -problem simultaneously). 2 More realistic formulation of the non-ideal MHD terms. Relating with a stochastic acceleration model of radio- $(\mathbf{3})$ emitting particle developed by Tanaka & Asano (2017).



 $v_{\rm Crab}$  ~ 1,500km/s because the MHD postshock flow is not spatially isobaric.

• Radial brightness profile would distinguish models, i.e., syn. luminosity strongly depends on  $\tau_{conv}$ ,  $au_{\rm diss}, \& r_{\rm TS}.$ 

3-8, Jun. 2019, Supernova Remnants II: An Odyssey in Space after Stellar death@Crete, Greece