Modeling PWNe and their magnetic fields Barbara Olmi

INAF, Osservatorio Astrofisico di Arcetri - Firenze - Italy Institute of Space Sciences (ICE, CSIC) - Barcellona - Spain

Pulsar Wind Nebulae

Broad band non-thermal spectrum





Gamma-rays: Inverse Compton scattering with local photon field



Adapted from Kennel & Coroniti 1984 [Del Zanna & Olmi 2017]



The central pulsar is both source of magnetic field and particles: it fills the remnant with a magnetized, relativistic and cold wind (mainly leptonic)

Adapted from Kennel & Coroniti 1984 [Del Zanna & Olmi 2017]



Interaction of the pulsar wind with the SNR induces the formation of a Termination shock

Adapted from Kennel & Coroniti 1984 [Del Zanna & Olmi 2017]



Adapted from Kennel & Coroniti 1984 [Del Zanna & Olmi 2017] The visible nebula corresponds to the shocked wind beyond the TS. The PWN bubble is formed by a hot plasma and intense magnetic field (50-200 µG)

17/27 static models of PWNe

[Rees & Gunn 1974, Kennel & Coroniti 1984, Emmering & Chevalier 1987, Begelman & Li 1992]



Assumptions:

- the cold isotropic relativistic PW terminates in a strong perpendicular shock
- the flow in the nebula is subsonic
- particle acceleration at the shock
- synchrotron losses beyond the shock

Main Free parameters:

- particle spectral indices
- wind Lorentz factor $\rightarrow \Gamma$
- wind magnetization $\rightarrow \sigma = B^2/(4\pi n m_e c^2 \Gamma^2)$

Predictions:

- positon of TS \rightarrow R_{TS} ~ R_N(V_N/c)^{1/2} ~ 0.1 pc
- Optical / X-ray spectrum [de Jager & Harding 1992, Atoyan & Aharonian 1996]
- size shrinkage with increasing energy

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- the cold isotropic relativistic PW terminates in a strong perpendicular shock
- the flow in the nebula is subsonic
- particle acceleration at the shock
- synchrotron losses beyond the shock

Main Free parameters:

- particle spectral indices
- wind Lorentz factor $-\Gamma$
- wind magnetization $\rightarrow \sigma = B^2/(4\pi n m_e c^2 \Gamma^2)$

From basic dynamics or radiation properties:

 $\Gamma \simeq 10^6$ $\sigma \simeq V_N/c \simeq 10^{-3}$

The sigma paradox

From pulsar theories $\rightarrow \sigma \sim 10^4 @ R_L$



From 1D PWNe models $\rightarrow \sigma \sim 10^{-3} \otimes R_{TS}$



Even if R_{TS}~10⁹ R_L dissipation is not sufficient to explain this discrepancy!

A deeper view in PWNe



[Crab Nebula - Chandra] Knots Arcs

[Vela Nebula - Chandra]

Counter jet

Equatorial torus

Polar jet

Outward moving wisps



Formation of the polar jets

Magnetic collimation in the relativistic PW is not efficient [Lyubarsky & Eichler 2001]: $\Gamma \gg 1 \rightarrow \rho \vec{E} + \vec{J} \times \vec{B} \sim 0$

Collimation must occur inside the nebula via hoop stresses. The energy flux is in the nebula is anisotropic! [Bogovalov & Khangoulian 2002, Lyubarsky 2002]

 $F \propto \sin^2(\theta)$

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TS is now oblate, with R_{eq} >> R_{pol}

2D numerical models of PWNe

[Michel 1973, Bogovalov 1999, Contopoulos + 1999, Coroniti 1990, Gruzinov 2004, Bogovalov & Khangoulian 2002, Lyubarsky 2002]

• Anisotropic distribution of the energy flux F(r,9):

0

Goals of 2D models: jet-torus

2D numerical models confirm the jet formation for values of magnetization σ≥0.01 [Komissarov & Lyubarsky 2003-2004, Del Zanna et al. 2004]





Goals of 27 models: variability of the inner nebula



Goals of 2D models: multi- λ wisps properties



Non-coincident locations and different outward velocities at different wavelengths explained with non uniform injection of emitting particles:



Limits of 2D models

 $\sigma_0 = 0.025$: Magnetic Field [μ G] Morphology of the magnetic field D Level of magnetization y [ly] 0 Total pressure ×1018 -6.0-6.22D simulations can only -5 work with $\sigma < 1$. -6.4-6.600 -6.8-5 × [ly] -7.0-7.2-200 -300 -1000 -7.4Averaged field [Olmi et al. 2015] underestimated: -7.6 $< B >_{SYM} \simeq 10^{-5} G$ $< B > OBS \simeq 10^{-4} G$ -7.8-8.00.0 0.5 1.0 1.5 $\times 10^{18}$ [Porth et al. 2014]

Limits of 2D models

This strongly affects many of the emitting properties:



Problems in reproducing high-energy spectrum: X-ray needs artificial steepening in order to compensate lower energy losses + IC overestimated

+ (14)



Moving to 3D: Crab models

3D models allow for a more complex structure of the magnetic field

✓ In 3D the magnetic dissipation is stronger (Kink instability) and $\sigma \ge 1$ can be reached!

Possible solution to σ -paradox?



BUT 3D simulations are demanding in terms of resources and time, data are huge



> 1 Million CPU hours (ran @ CINECA with >2000 CPUs)







age t=250 yr

[Olmi et al. 2016]

Not everything is solved...

The average field is again too low: $\sigma=1$ not sufficient?



Same problems with the spectrum!

3D vs 2D

global dynamics completely different

inner dynamics comparable



[Porth et al. 2013, 2014]

Evolved PWNe



Large fraction of all the pulsars born with high kick velocity (10%-50%)



fated to escape the SNR on timescales << than typical pulsar ages (~10⁶ yr)

Bow shock nebulae



Bow shock nebulae: cometary shape









PSR J1509-5850

X-ray Radio 15" Radio

Mouse PWN [Yusef-Zadeh & Bally 1987, Yusef-Zadeh & Gaensler 2005, Klinger et al. 2018]

PSR J1509-5850 [Hui & Becker 2007, Klinger et al. 2016]

Bow shock nebulae: puzzling outflows and halos





Geminga [Posselt et al. 2017]



Extended TeV halo [Abeysekara et al. 2017]

X-ray





Lighthouse nebula [Pavan et al. 2016]

Guitar nebula [Cordes et al. 1993, Wong et al. 2003]



G327 [Temim et al. 2009]

PSR J1509-5850 [Klinger et al. 2016]

2D MHD Models of BSPWNE

Relativistic MHD axisymmetric simulations [Bucciantini et al. 2005] were able to account for the formation of the bow shock, TS deformation, generation of the tail.



Tail with cylindric shape with constant area.

3D MHD Models of BSPWNe

How to model the pulsar wind:

Inclination of the spin-axis and pulsar speed:



Spin-axis aligned with pulsar motion $\Phi_M=0^\circ$

Φ_M=45°

Φ_M=90°

Anisotropy in the energy flux: $F(\psi) \propto 1 + \alpha \sin^2 \psi$ ψ colatitude from the spin-axis



Wind magnetization: 0.01 $\lesssim \sigma$

Effects of magnetic field geometries on the BS [Barkov & Lyutikov 2018]



Anisotropy of ISM density only slightly affect the FS morphology [Toropina 2018, Barkov & Lyutikov 2018]



Simulations supported by MoU INAF-CINECA class A projects

Dynamics of the tail [Olmi & Bucciantini 2019]



Maps of the magnetic field for different geometries and magnetizations

Different field geometries (inclinations)

Shape of the contact discontinuity

Magnetization + isotropy/anisotropy

Dynamics in the tail

Development of turbulence [Olmi & Bucciantini 2019]



Low magnetization high level of turbulence, chaotic flow

High magnetization

low level of turbulence, flow almost laminar

Development of turbulence [Olmi & Bucciantini 2019]



Development of Eurbulence [Olmi & Bucciantini 2019]



Development of turbulence



Maps of the magnetization

HIGH TURBULENCE with complete loss

of injection information



LOW TURBULENCE with injection information maintained and laminar flow

[Olmi & Bucciantini 2019]

Consequences for emission



Preliminary from Olmi & Bucciantini 2, in prep.

Escape of particles



Escape of particles

Particles in the polar flow are confined by the currents in the magnetopause layer.



From the reconnection point particles stream out along the ISM magnetic field, forming jets (here symmetric since the wind is symmetric).



Preliminary from Olmi & Bucciantini 3, in prep.

Conclusions

Take home message

Magnetic field can be realistically modeled ONLY in 3D

2D MHD simulations can only reproduce the inner nebula properties, where deviations from 3D are small.

A realistic morphology matching with the overall emission properties based on 2D model is more safe with HD simulations:



G21.5-0.9 simulations, Olmi & Torres in prep.



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Thank you!!!

