

A vibrant, multi-colored image of the Tycho supernova remnant, showing a complex, irregular shape with a mix of yellow, green, and blue hues against a black background.

# HOW COSMIC RAYS SHAPE SUPERNOVA REMNANTS

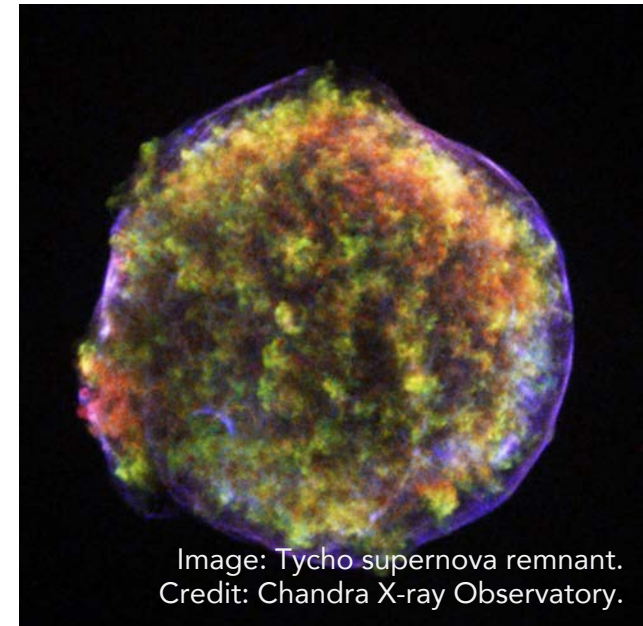
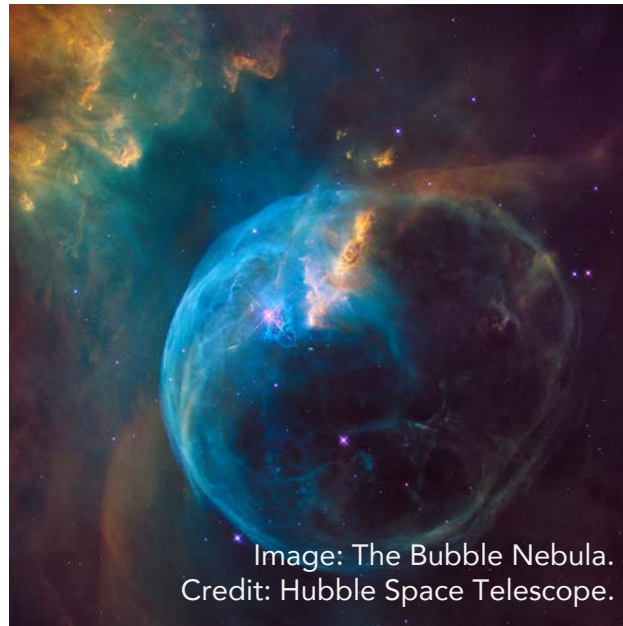
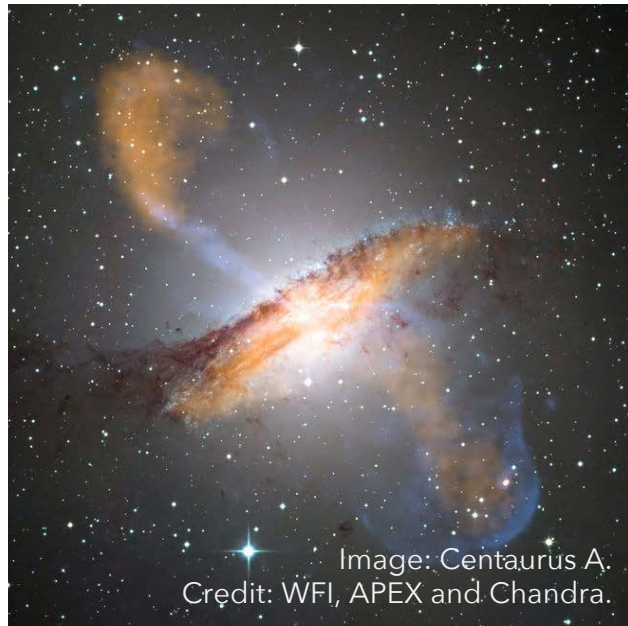
REBECCA DIESING  
SUPERNOVA REMNANTS II

Image: Tycho supernova remnant. Credit: Chandra X-ray Observatory.

# MOTIVATION

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Feedback from active galactic nuclei, stellar winds, and supernovae plays a crucial role in galaxy formation.



These objects inject energy and momentum into the interstellar medium.

# MOTIVATION

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To produce galaxies consistent with observations, models typically boost the momentum yield calculated for single supernova remnants (SNRs) by a factor of 3-5.

# THE ROLE OF COSMIC RAYS

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Cosmic rays (CRs) are produced at SNR blast waves via diffusive shock acceleration. Their effect is two-fold:

1. Acting as a relativistic fluid, CRs suffer less adiabatic loss than thermal gas; at late times they dominate the internal pressure.

# THE ROLE OF COSMIC RAYS

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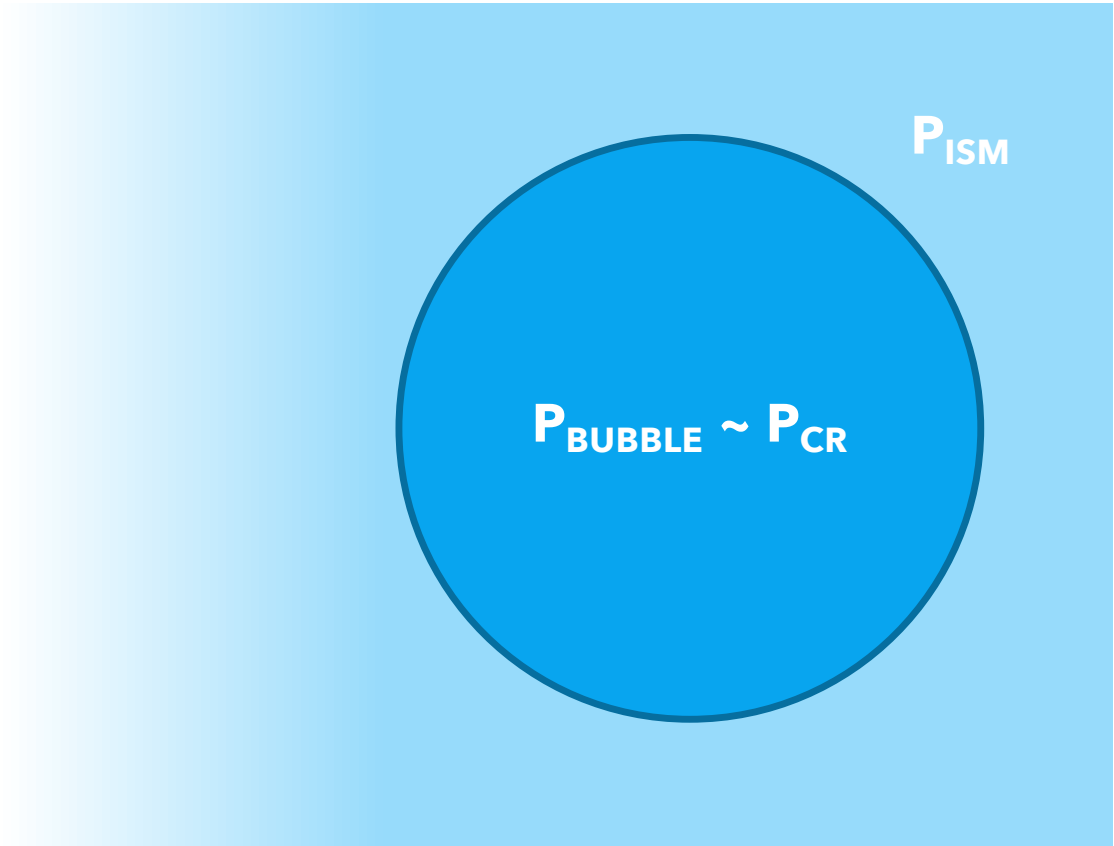
Cosmic rays (CRs) are produced at SNR blast waves via diffusive shock acceleration. Their effect is two-fold:

1. Acting as a relativistic fluid, CRs suffer less adiabatic loss than thermal gas; at late times they dominate the internal pressure.
2. CR energy is not radiated at late times, but rather continues to support expansion.

# A SIMPLE ESTIMATE

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After the thermal gas radiates away its energy, cosmic rays push the remnant outward.

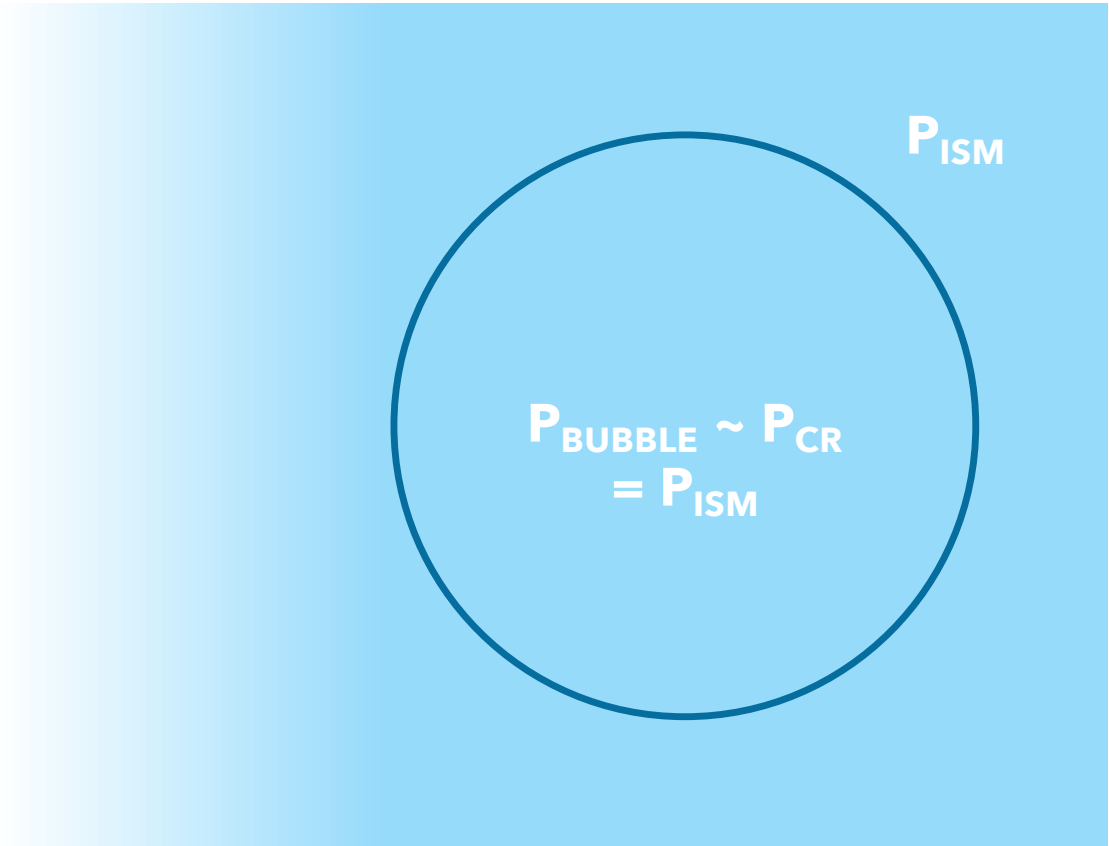


# A SIMPLE ESTIMATE

After the thermal gas radiates away its energy, cosmic rays push the remnant outward.

Momentum becomes constant when pressures equilibrate:

CR acceleration efficiency  $\nearrow$   $\frac{\xi E_{\text{SN}}}{V} = n_{\text{ISM}} kT$





# A SIMPLE ESTIMATE

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Assume ejecta mass is negligible and estimate the final SNR mass:

$$M_f \approx n_{\text{ISM}} \mu m_p V = \frac{\xi E_{\text{SN}} \mu m_p}{kT}$$

Assume final SNR velocity is of order the ISM sound speed:

$$v_f = c_s = \sqrt{\frac{\gamma kT}{\mu m_p}}$$



# A SIMPLE ESTIMATE

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Momentum deposited is simply the final mass times the final velocity:

$$p_{\text{dep}} = M_f v_f \approx \xi E_{\text{SN}} \sqrt{\frac{\gamma \mu m_p}{kT}}$$

Assuming  $\gamma = 5/3$  and  $\mu = 1.4$ :

$$p_{\text{dep}} \approx 9.44 \times 10^5 \frac{\xi}{0.1} \frac{E_{\text{SN}}}{10^{51} \text{erg}} \left( \frac{T}{8000 \text{K}} \right)^{-1/2} M_{\odot} \text{km s}^{-1}$$

# A SIMPLE ESTIMATE

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WITH COSMIC RAYS

$$p_{\text{dep}} \approx 9.44 \times 10^5 \frac{\xi}{0.1} \frac{E_{\text{SN}}}{10^{51} \text{erg}} \left( \frac{T}{8000 \text{K}} \right)^{-1/2} M_{\odot} \text{km s}^{-1}$$

WITHOUT COSMIC RAYS

$$p_{\text{dep}} \approx 2.8 \times 10^5 M_{\odot} \text{km s}^{-1} \leftarrow \text{Factor of } \sim 3 \text{ lower than SNRs with CRs}$$

# A MORE DETAILED CALCULATION

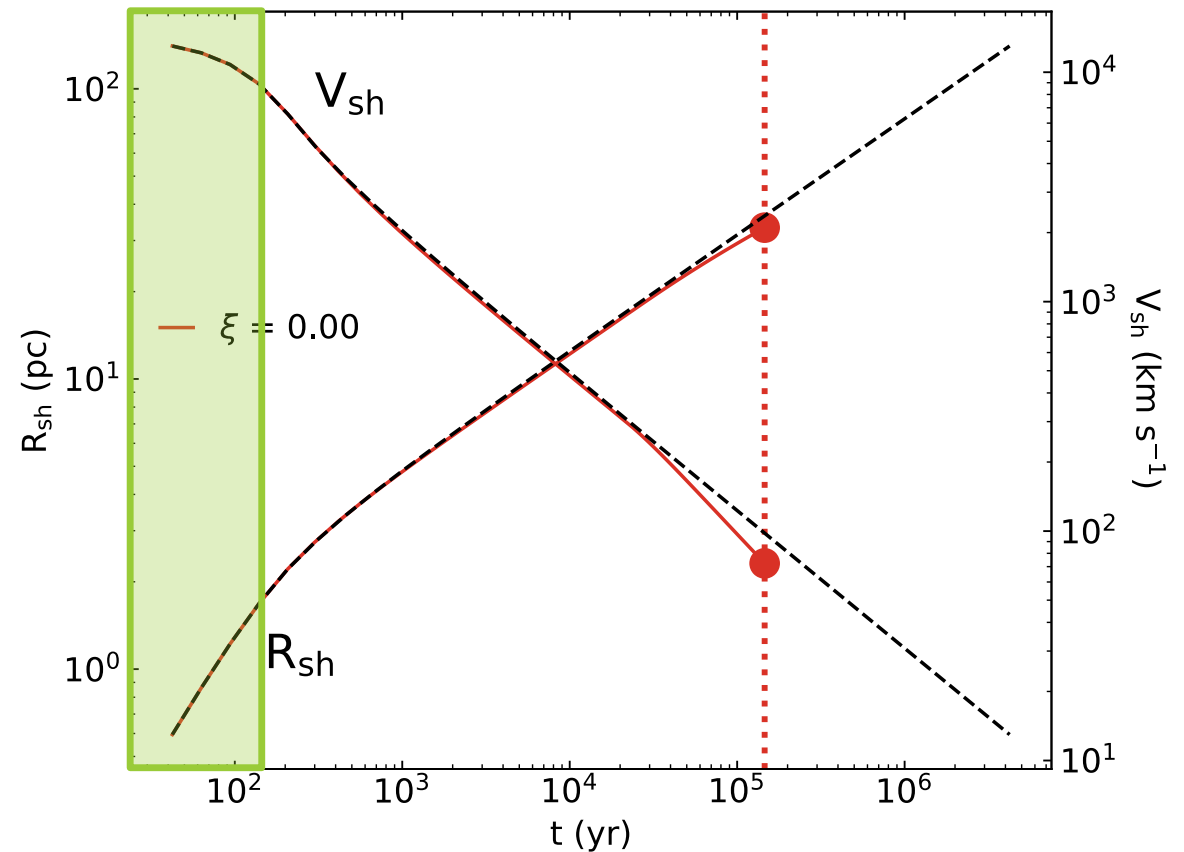
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# 1. EJECTA-DOMINATED STAGE

Energy is conserved and swept up mass  $<$  ejecta mass.

CR contribution to expansion is negligible.

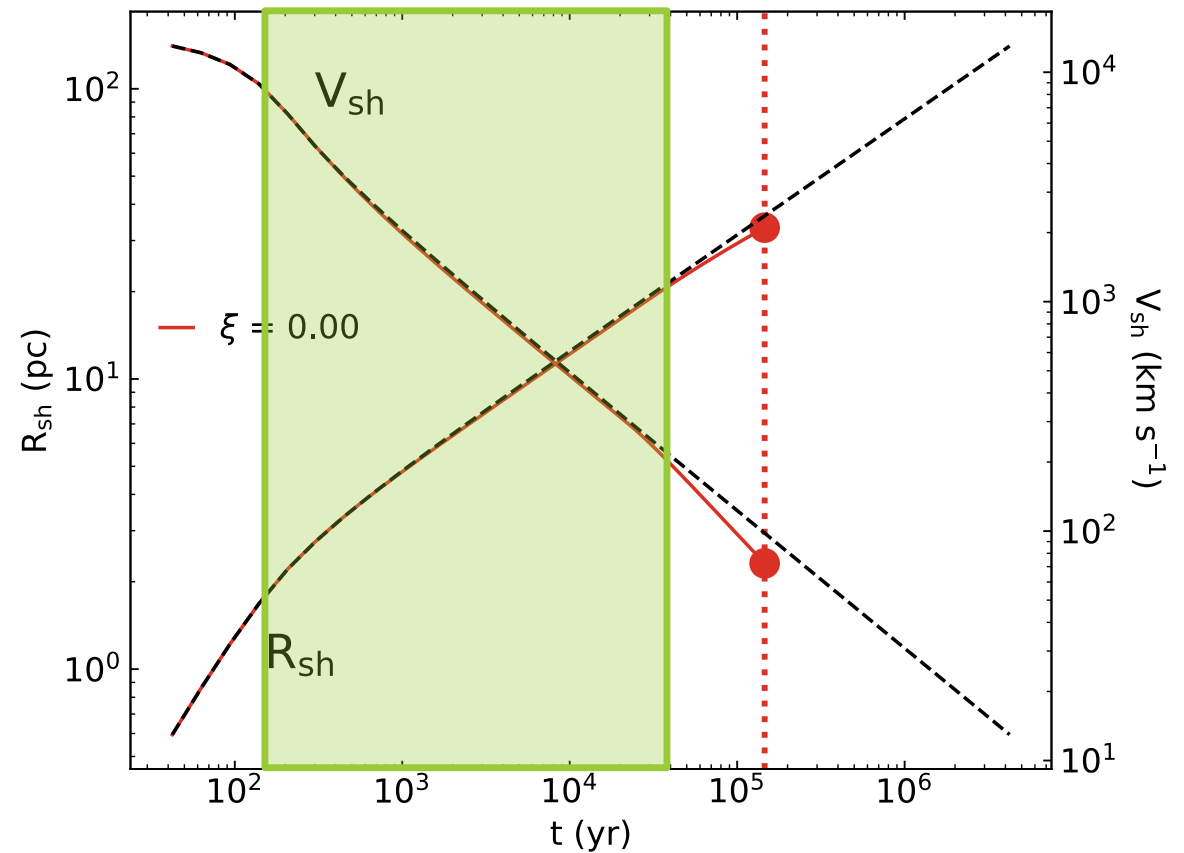
Use analytic approximation from Truelove & McKee 1999 (dashed line).



## 2. SEDOV-TAYLOR STAGE

Energy is conserved and swept up mass  $>$  ejecta mass.

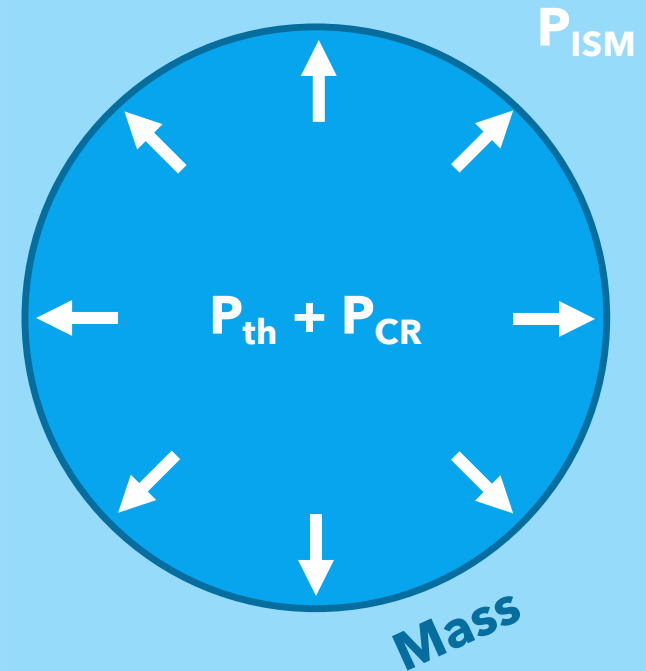
Use thin-shell approximation.



# THIN SHELL APPROXIMATION

Assume mass is confined to a thin shell.

Shell expands due to pressure inside the hot bubble.



# THIN SHELL APPROXIMATION

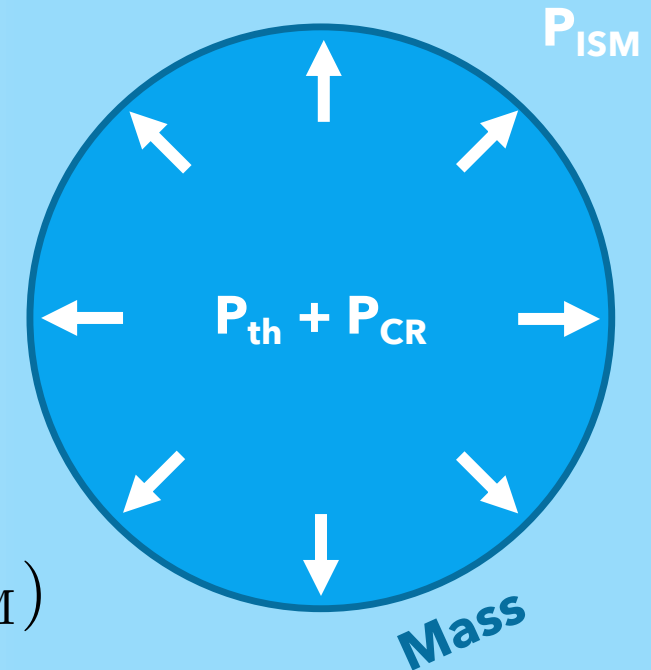
Enforce conservation laws:

$$M(R_{sh}) = M_{ej} + \frac{4\pi}{3} R_{sh}^3 \rho_{ISM}$$

Effective  
adiabatic index  
depends on  $\xi$

$$E = E_{SN}$$

$$\frac{2}{\gamma_{\text{eff}} + 1} \frac{d(Mv_{sh})}{dt} = 4\pi R_{sh}^2 (P_{th} + P_{CR} - P_{ISM})$$

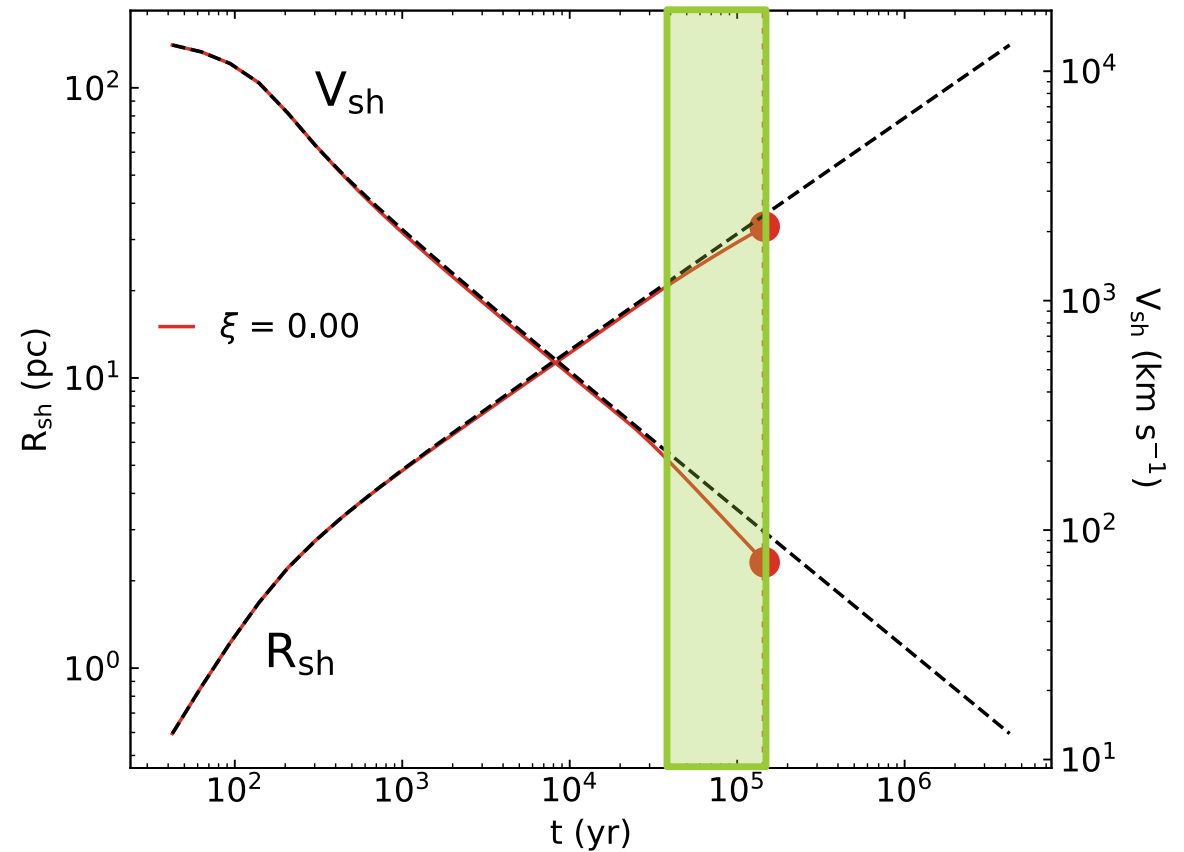




# 3. RADIATIVE STAGE

$T \lesssim 10^6 \text{ K} \rightarrow$  thermal gas radiates away energy.

Use modified thin-shell approximation.



# MODIFIED THIN SHELL APPROXIMATION

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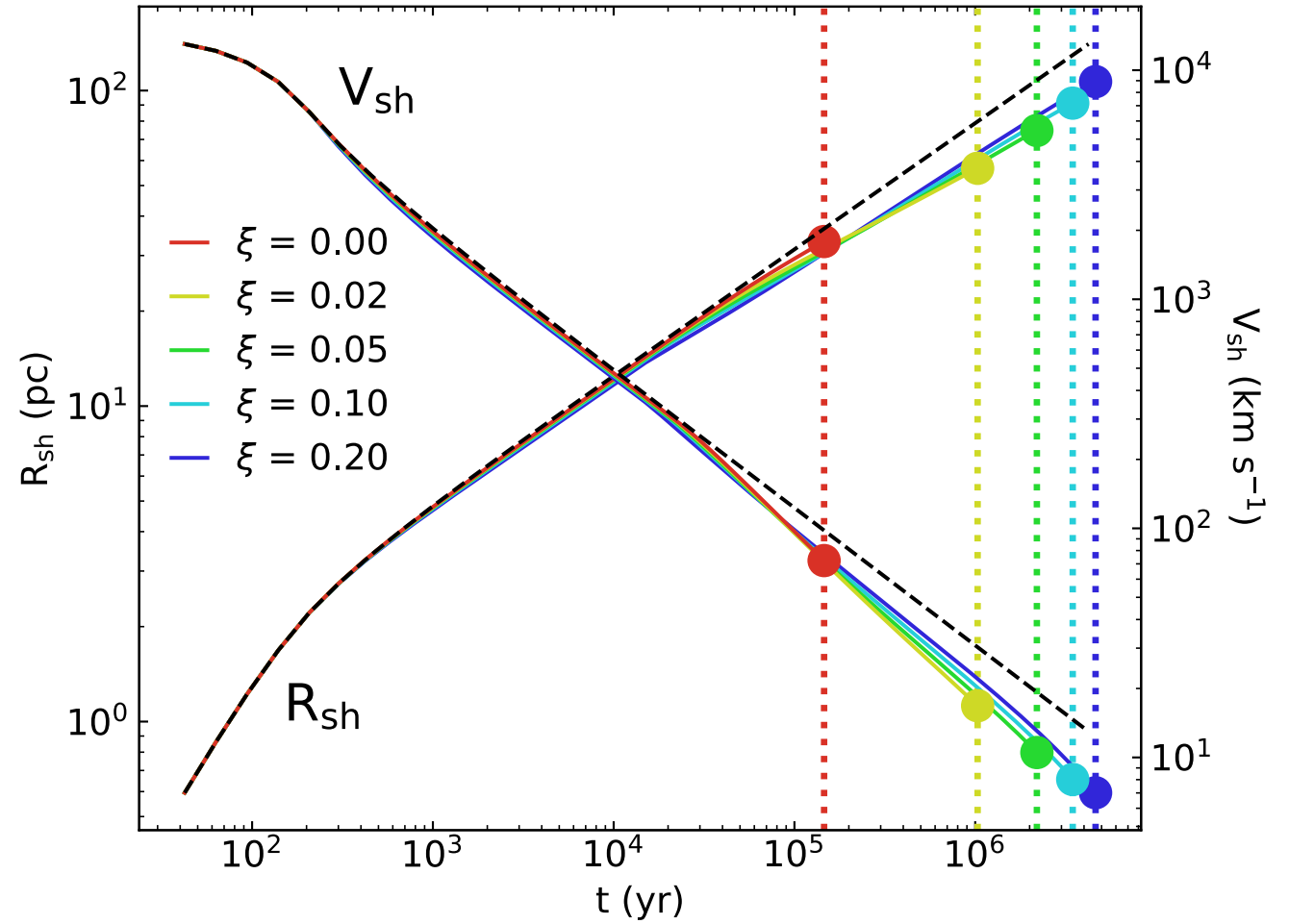
Enforce new conservation laws:

$$M(R_{sh}) = M_{ej} + \frac{4\pi}{3} R_{sh}^3 \rho_{\text{ISM}}$$

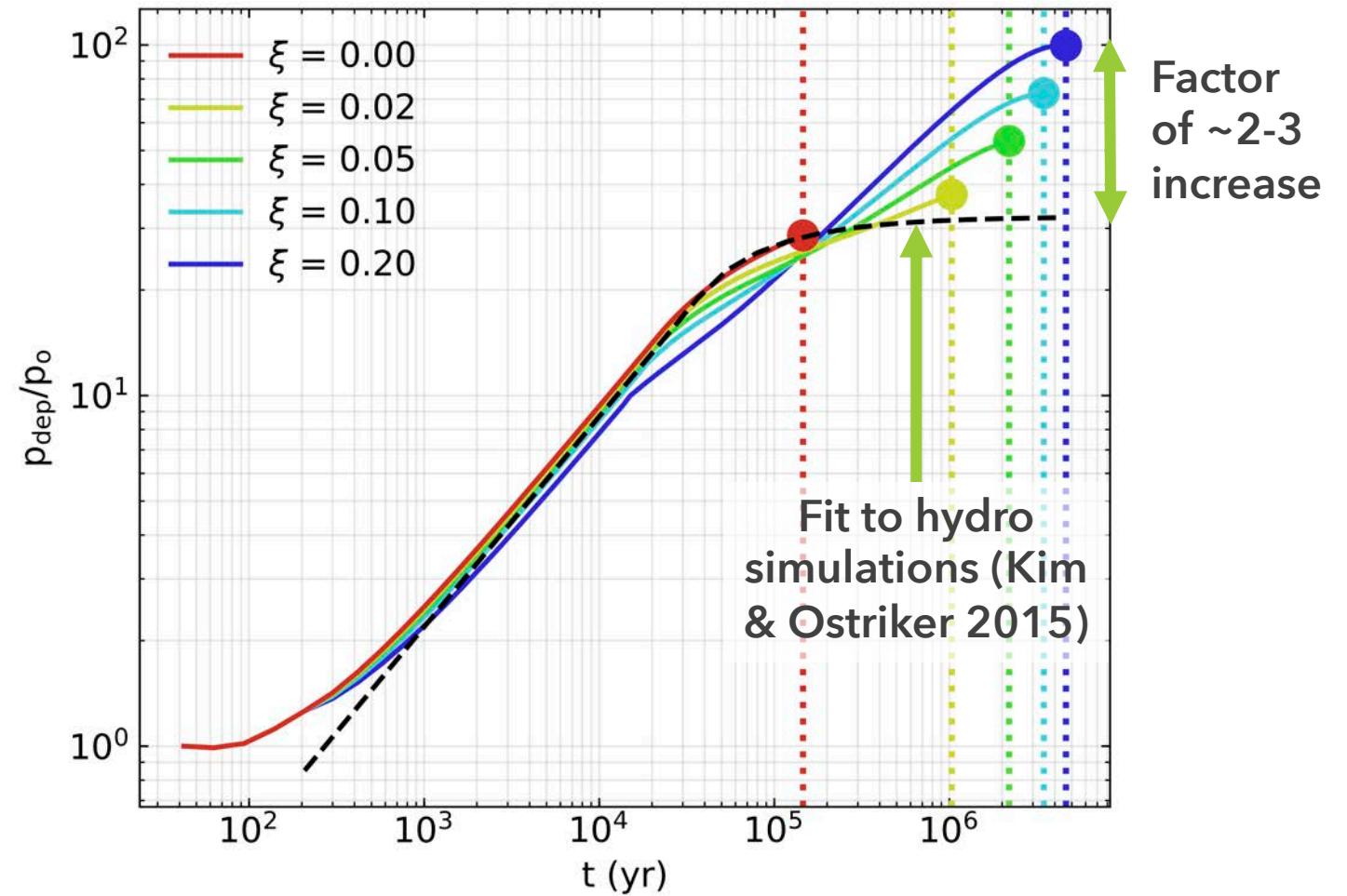
$$E(R_{sh}) = E_{SN} - \int_{t_{\text{rad}}}^t d\tau 2\pi \chi_{th} \rho v_{sh}^3 R_{sh}^2$$

$$\frac{2}{\gamma_{\text{eff}} + 1} \frac{d(Mv_{sh})}{dt} = 4\pi R_{sh}^2 (P_{th} + P_{\text{CR}} - P_{\text{ISM}})$$

# RADIUS & VELOCITY EVOLUTION

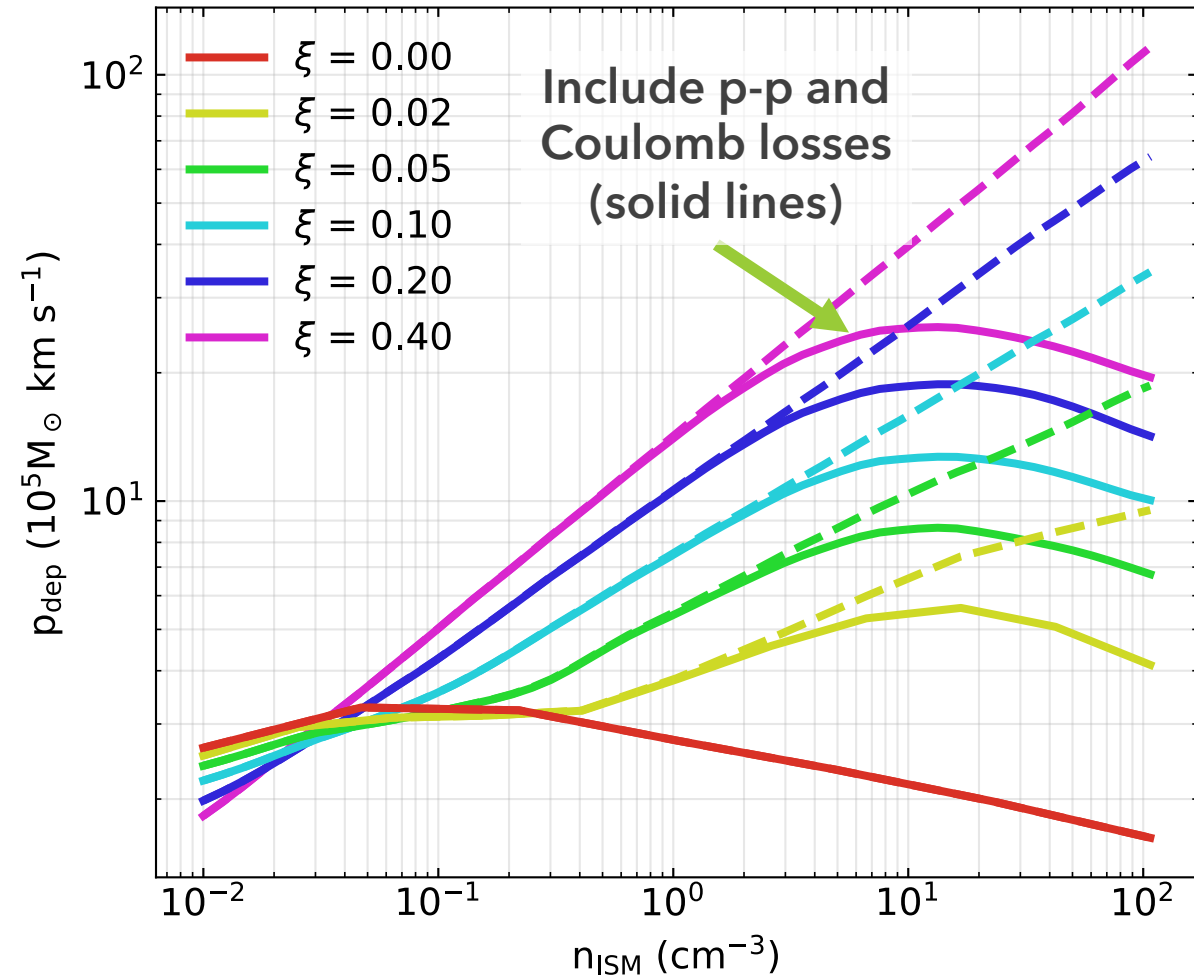


# MOMENTUM EVOLUTION



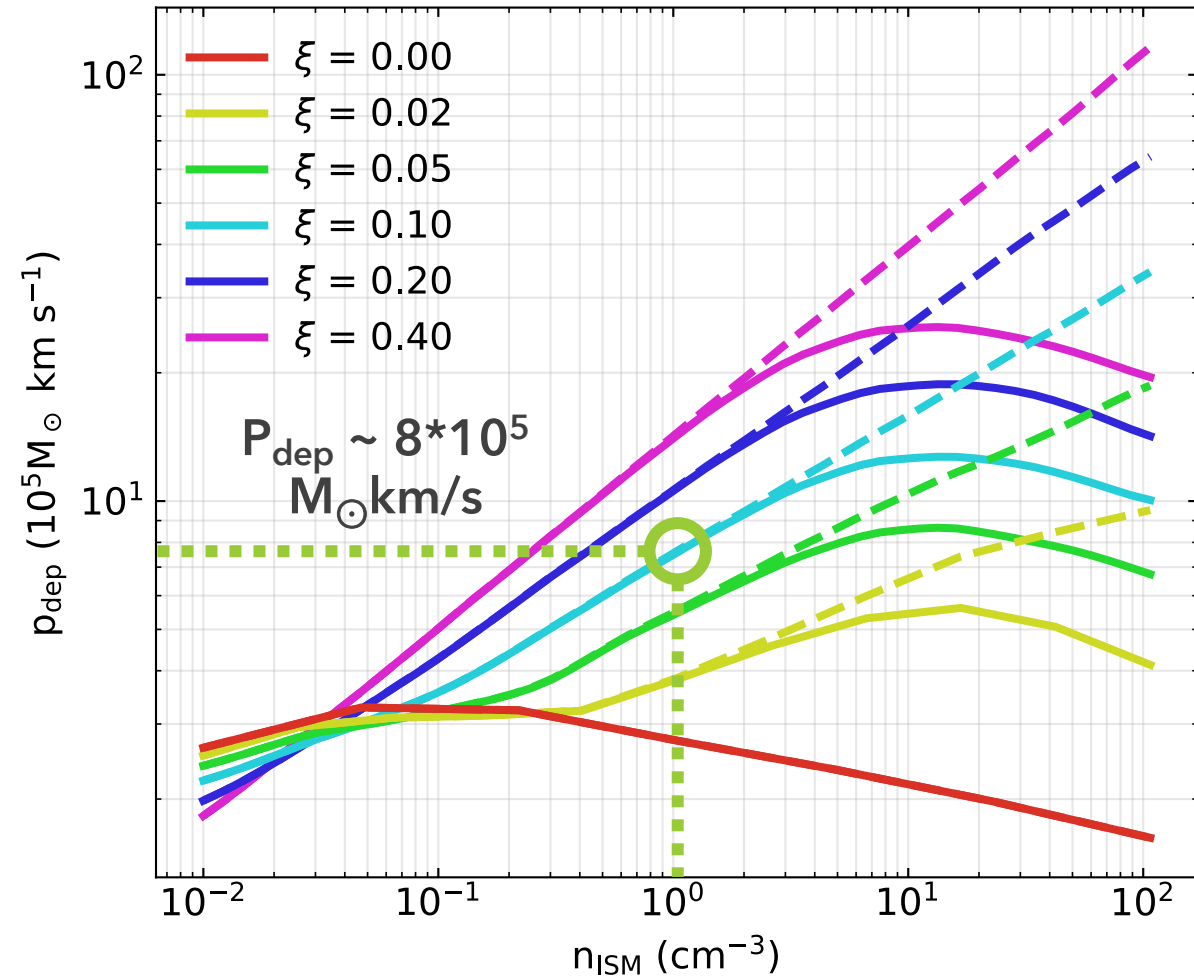
# EFFECT OF ISM DENSITY

(keeping pressure constant)



# EFFECT OF ISM DENSITY

(keeping pressure constant)



# CONCLUSIONS

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1. **This work represents the first semi-analytical calculation—in the thin-shell approximation limit—of SNR evolution through the Sedov-Taylor and radiative stages.**  
Our approach accurately reproduces the main features of hydro simulations (w/o CRs).
1. **The presence of cosmic rays prolongs the expansion of SNRs, resulting in 2-3 times more momentum deposition in the ISM.**  
This boost becomes more pronounced in dense environments.
2. **Galaxy formation simulations that include supernova feedback via sub-grid models should account for the additional contribution of the CRs accelerated at SNR shocks.**

For more information, please see Diesing & Caprioli 2018.